

On the 'DSF' and the 'Dreaded Surface Effect'

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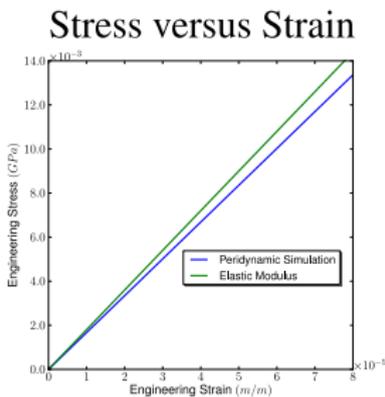


What is the DSF?

For Isotropic-Ordinary Materials

It *MIGHT* be a Practical Solution to the Following Problem

Axial Displacement



Stored Elastic Energy



The following related aspects contribute to the above mismatch.

- Geometric surface effects
- Nonlocal model (dilatation on surface) and model properties
- Discretization error

This talk is about working towards a simple/practical solution.

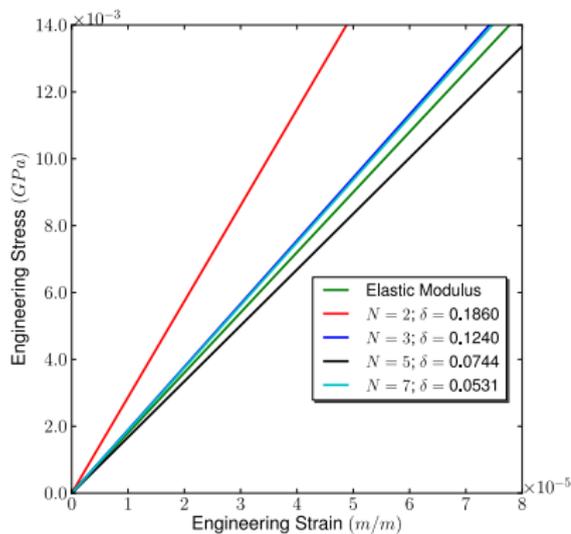


Thanks for the Help and Support

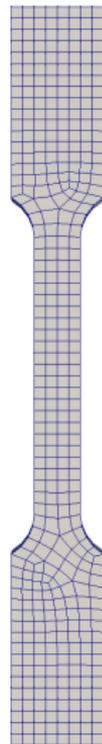
- **Dave Littlewood** – generally helpful with running *Peridigm* and especially for providing the demonstration problem
- **Stewart Silling** – discussion about the 'DSF' and approach presented here today.
- **Mike Parks** – *Peridigm* support

Mesh Refinement Study

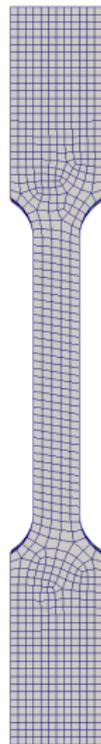
Horizon is tied to mesh element size h : $\delta = 3h$



$N = 2$



$N = 3$



$N = 7$



inches



Review of Isotropic-Ordinary Materials

Scalar Force State $t = \frac{3K\theta}{m}\omega|\xi| + \alpha\omega e^d$

$$e = |Y| - |\xi| \quad \theta = \frac{3}{m}(\omega|\xi|) \bullet e \quad e^d = |Y| - |\xi| - \frac{\theta}{3}|\xi|$$

Substituting the above relations

$$t = \left(\frac{3K}{m} - \frac{\alpha}{3} \right) \theta \omega |\xi| + \alpha \omega e$$

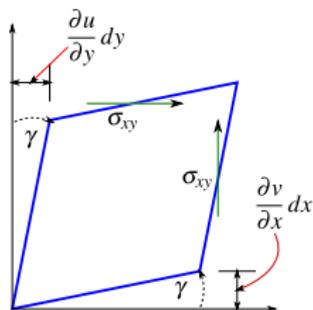
If we use $\alpha = \frac{15\mu}{m}$ (trouble starts here)

$$t = \left(\frac{9K - 15\mu}{3m} \right) \theta \omega |\xi| + \alpha \omega e$$

We need a 'discretization' appropriate and 'geometric' appropriate α



infinitesimal shear deformation St. Venant-Kirchhoff Material



Displacements

$$\begin{aligned}u &= \gamma y \\v &= \gamma x \\w &= 0\end{aligned}$$

Deformation Gradient F

$$[F] = \begin{bmatrix} 1 & \gamma & 0 \\ \gamma & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Green Lagrange-Strain E

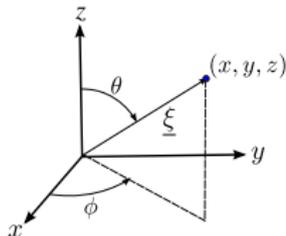
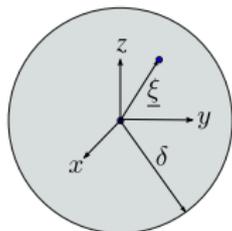
$$\begin{aligned}2[E] &= [F]^t[F] - I \\ &= \begin{bmatrix} \gamma^2 & 2\gamma & 0 \\ 2\gamma & \gamma^2 & 0 \\ 0 & 0 & 0 \end{bmatrix}\end{aligned}$$

Strain energy density $W_L = \frac{1}{2}\lambda[tr(E)]^2 + \mu E : E = \frac{1}{2}\lambda\gamma^4 + 2\mu\gamma^2$



infinitesimal shear deformation Ordinary Isotropic Peridynamic Material

On a sphere



Deformation State Y

$$Y(\gamma) = \sqrt{z^2 + (y + \gamma x)^2 + (x + \gamma y)^2}$$

Linearize deformation state

$$Y(\gamma) \approx \dots + \mathcal{O}(\gamma^2)$$

Form extension state $e = Y - |\xi|$

$$e = \gamma |\xi| \sin^2(\theta) \sin(2\phi)$$

Verify dilatation is zero

$$\frac{3}{m} |\xi| \bullet e = 0$$

Weighted volume on a sphere

$$m = \frac{4\pi\delta^5}{5}$$

$$\text{Energy density } W = \frac{1}{2} \alpha e^d \bullet e^d = 2\alpha \frac{m\gamma^2}{15}$$

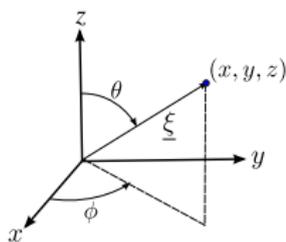
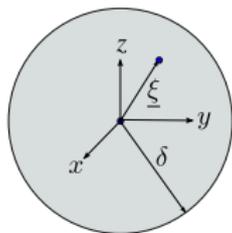


Elastic energy density for *infinitesimal* shear deformation

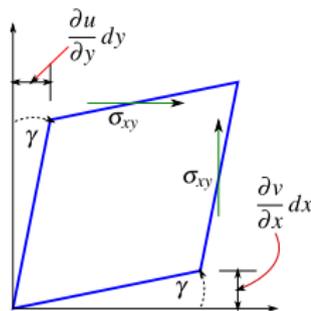
Review Stewart's Approach

Equating nonlocal W with local W_L

On a sphere



$$W = \frac{1}{2} \alpha e^d \bullet e^d = 2\alpha \frac{m\gamma^2}{15}$$



$$W_L \approx \mu \gamma^2$$

$$W_L = W \implies \alpha = \frac{15\mu}{m}$$

Observations

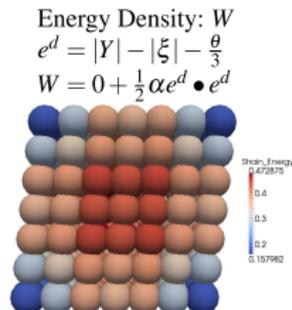
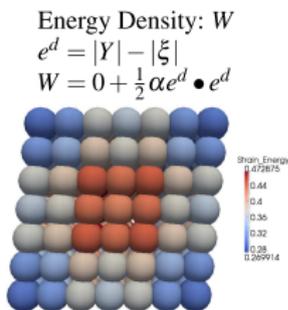
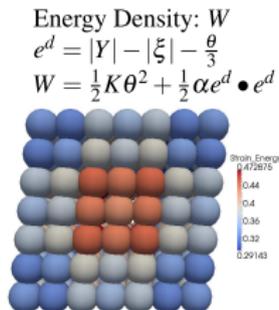
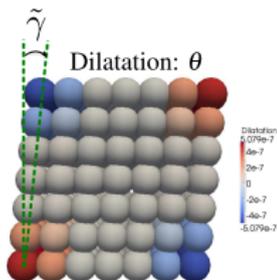
- W_L is independent of position, i.e., *constant*
- W depends upon neighborhood through integral.



Elastic energy density

Consider simple shear: $u = \tilde{\gamma}y$; $v = 0$; $w = 0$; $\gamma = \frac{\tilde{\gamma}}{2}$; $W_L = \frac{1}{2}\mu\tilde{\gamma}^2$

$\mu = 6.923 \times 10^{11}$; $K = 1.5 \times 10^{12}$; $\tilde{\gamma} = 1.0 \times 10^{-6}$; $W_L \approx .34615$



Motivated by Stewart's Approach

Create a Surface Correction Factor: DSF

Reconsider elastic energy density for shear at a point

$$W = \frac{1}{2} \tilde{\alpha} ||e^d||_h^2 = W_L = 2\mu \gamma^2$$

Divide both sides by $||e^d||^2 = \frac{4m\gamma^2}{15}$, where $m = \frac{4\pi\delta^5}{5}$

$$\begin{aligned} \tilde{\alpha} &= \frac{||e^d||^2}{||e^d||_h^2} \frac{15\mu}{m} \\ &= \frac{4m}{15} \frac{\gamma^2}{||e^d||_h^2} \frac{15\mu}{m} \\ &= DSF \times \alpha \end{aligned}$$

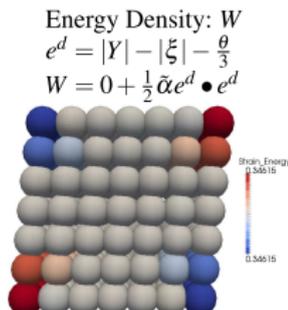
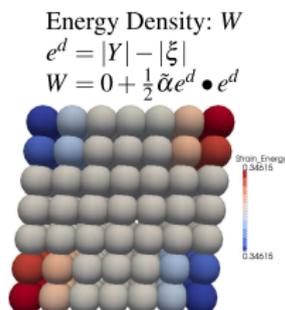
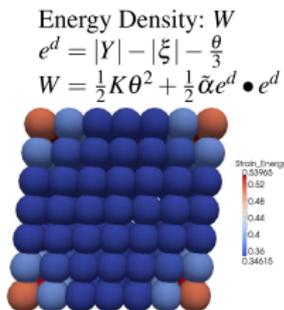
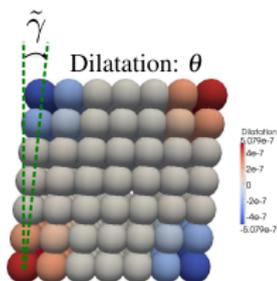


Reconsider Simple Shear Example Using *DSF*

Consider simple shear: $u = \tilde{\gamma}y$; $v = 0$; $w = 0$; $\gamma = \frac{\tilde{\gamma}}{2}$; $W_L = \frac{1}{2}\mu\tilde{\gamma}^2$

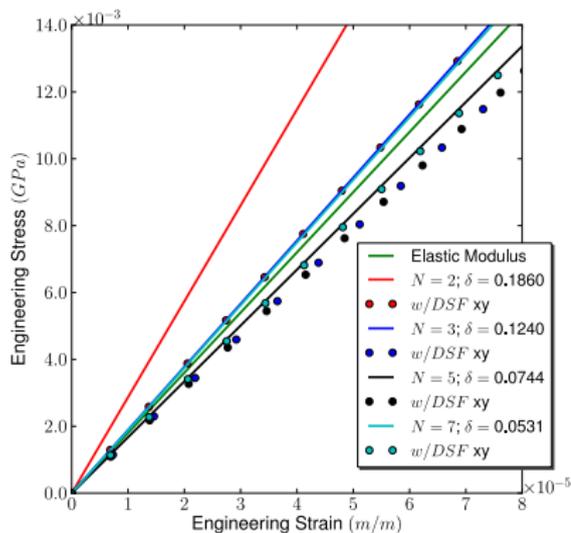
$\mu = 6.923 \times 10^{11}$; $K = 1.5 \times 10^{12}$; $\tilde{\gamma} = 1.0 \times 10^{-6}$; $W_L \approx .34615$

$$\alpha = \frac{15\mu}{m}; \quad \tilde{\alpha} = DSF \times \alpha$$



Mesh Refinement Study w/DSF

Horizon is tied to mesh element size h : $\delta = 3h$



$N = 2$



$N = 5$



Algorithm for computing the *DSF*

For all points *XP* in the neighborhood of *X*

Apply pure shear deformation

```
double dx = XP[0]-X[0];
double dy = XP[1]-X[1];
double dz = XP[2]-X[2];
double zx(0.0), xy(0.0), yz(0.0);
double xz(0.0), yx(0.0), zy(0.0);
switch(mode) {
case ZX:
    zx = gamma * dx;
    xz = gamma * dz;
    break;
case XY:
    xy = gamma * dy;
    yx = gamma * dx;
    break;
case YZ:
    yz = gamma * dz;
    zy = gamma * dy;
    break;
}
YP[0] = XP[0] + xy + xz;
YP[1] = XP[1] + yz + yx;
YP[2] = XP[2] + zx + zy;
```

Compute *DSF*

```
double scf, max_dsf, DSF;
int mode = XY;
Y[0] = X[0]; Y[1] = X[1]; Y[2] = X[2];
apply_pure_shear(X, Y, XP, YP, mode, gamma);
scf=ed_squared(neigh[X], X, Y, mode);
max_dsf=scf;

mode = ZX;
Y[0] = X[0]; Y[1] = X[1]; Y[2] = X[2];
apply_pure_shear(X, Y, XP, YP, mode, gamma);
scf=ed_squared(neigh[X], X, Y, mode);
if(scf>max_dsf) max_dsf=scf;

mode = YZ;
Y[0] = X[0]; Y[1] = X[1]; Y[2] = X[2];
apply_pure_shear(X, Y, XP, YP, mode, gamma);
scf=ed_squared(neigh[X], X, Y, mode);
if(scf>max_dsf) max_dsf=scf;

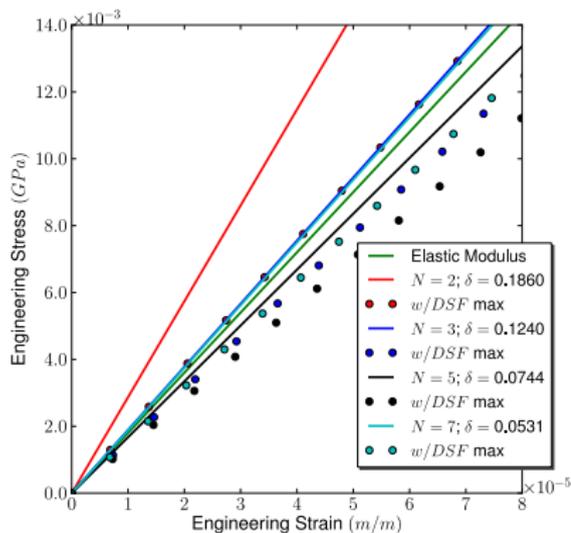
scf=max_dsf;
DSF = 4.0 * gamma * gamma * m / scf /15.0;
```

Computation on right uses *max*



Mesh Refinement Study w/DSF

Horizon is tied to mesh element size h : $\delta = 3h$



$N = 2$

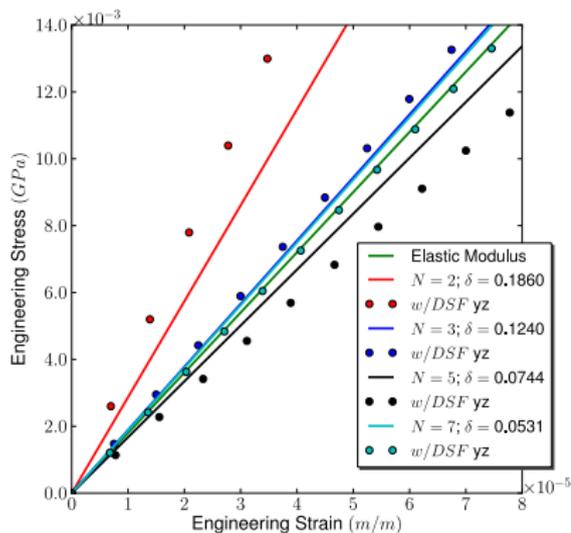


$N = 5$

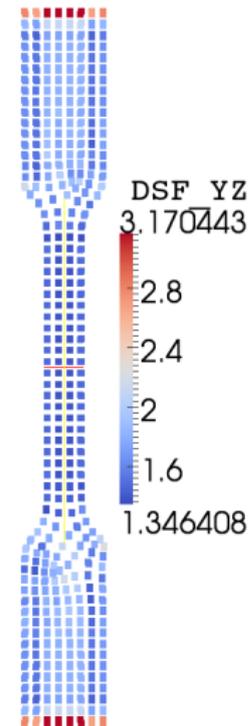


Mesh Refinement Study w/DSF

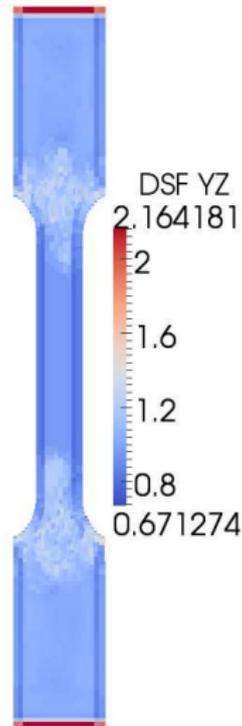
Horizon is tied to mesh element size h : $\delta = 3h$



$N = 2$

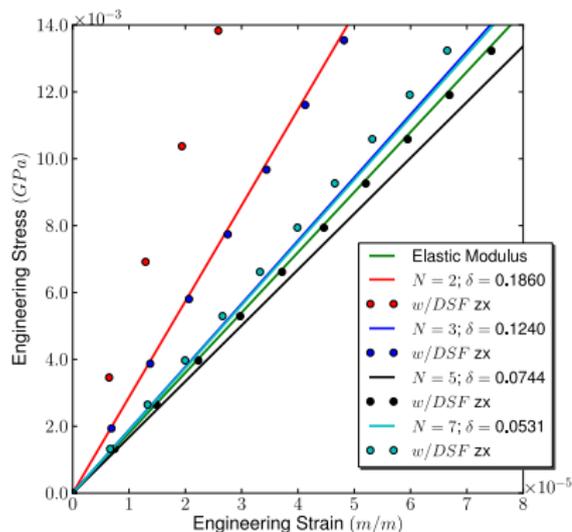


$N = 5$

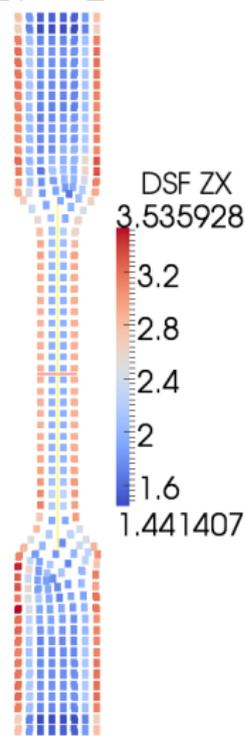


Mesh Refinement Study w/DSF

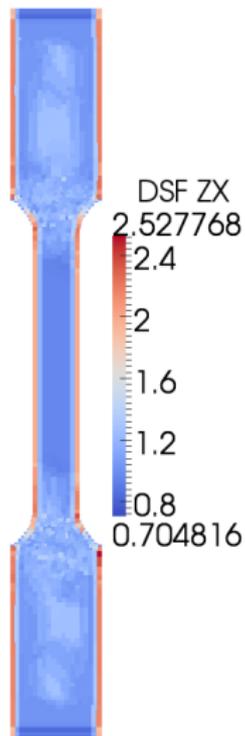
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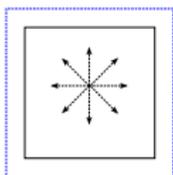


$N = 2$



$N = 5$





Displacements

$$u = \Delta x$$

$$v = \Delta y$$

$$w = \Delta z$$

Deformation Gradient F

$$[F] = \begin{bmatrix} 1 + \Delta & 0 & 0 \\ 0 & 1 + \Delta & 0 \\ 0 & 0 & 1 + \Delta \end{bmatrix}$$

Green Lagrange-Strain E

$$\begin{aligned} 2[E] &= [F]^t[F] - I \\ &= \begin{bmatrix} \Delta^2 + 2\Delta & 0 & 0 \\ 0 & \Delta^2 + 2\Delta & 0 \\ 0 & 0 & \Delta^2 + 2\Delta \end{bmatrix} \end{aligned}$$

Infinitesimal assumption: retain only Δ^2 terms

Strain energy density $W_L = \frac{1}{2}\lambda[tr(E)]^2 + \mu E : E = \frac{1}{2}(9\lambda + 6\mu)\Delta^2$

Closing Remarks and Summary

Expectation for isotropic ordinary materials

- Recover *Youngs Modulus* on a relatively **coarse** mesh

Using an energy based approach

- Developed shear modulus correction that shows promise
- Speculate that practical/ideal 'DSF' exists
- Question: Is the ideal 'DSF' *anisotropic* at the surface?

Identified issue

- Nonzero dilatation on surface under conditions of pure shear

Surface effect diminishes with mesh refinement

- Convergence is slow but apparent
- On *tensile test*, rate of convergence is inadequate

