

# Lagrange-Multiplier-Based Partitioned Method for Ocean- Atmosphere Coupling



VIII International Conference on Coupled Problems  
in Science and Engineering

Barcelona, Spain June 3-5, 2019

Kara Peterson, Pavel Bochev, Paul Kuberry, Nat Trask

Sandia National Laboratories

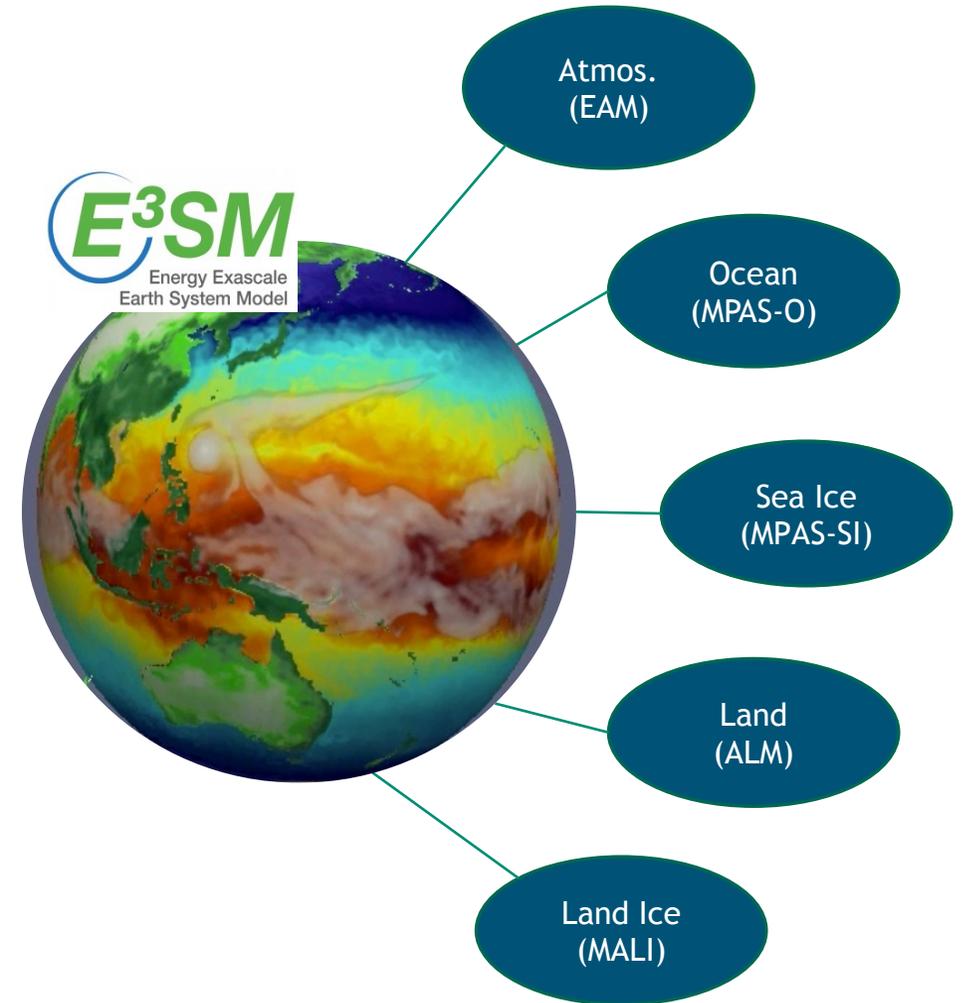


CANGA  
Coupling Approaches  
for Next Generation  
Architectures



## Earth System Model Coupling

- ESMS include multiple components for the ocean, atmosphere, ice, etc.
- Coupled problem is a complex multi-physics, multiscale problem
- Monolithic solutions of the coupled problem not computationally feasible
- Need stable and accurate methods for partitioned solves
- *Challenges:*
  - Non-conforming grids
  - Independent discretizations
  - Flux conservation and property preservation
  - Stability over long integration times

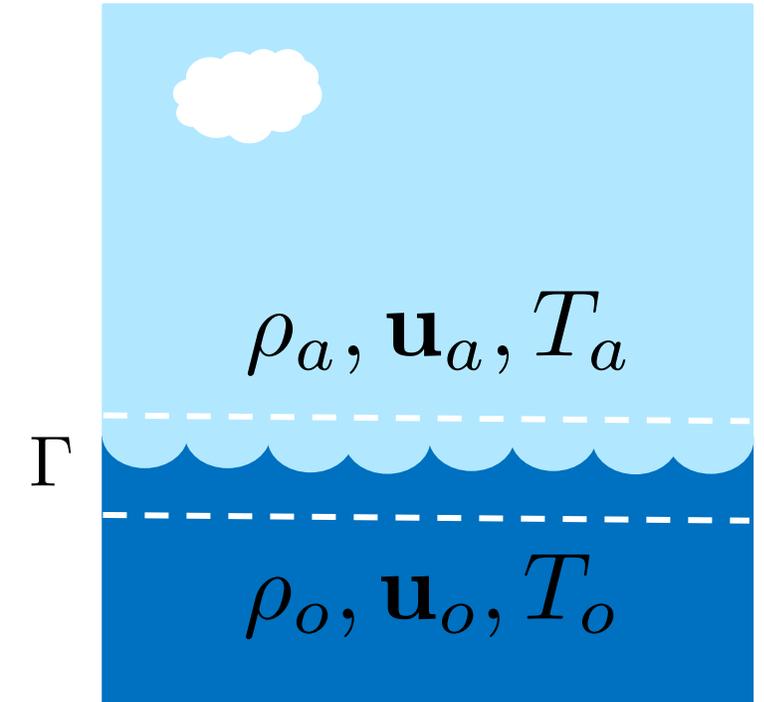


- Consider partial differential equations for atmosphere and ocean circulation with state variables **velocity** and **temperature**
- Ocean-atmosphere fluxes are defined by a parameterization of the surface layers: “bulk” formulation
- Coupling conditions

$$\rho_a \nu_a \frac{\partial \mathbf{u}_a}{\partial z} = \rho_o \nu_o \frac{\partial \mathbf{u}_o}{\partial z} = \boldsymbol{\tau} \quad \text{on } \Gamma$$

$$\rho_a K_a \frac{\partial T_a}{\partial z} = \rho_o K_o \frac{\partial T_o}{\partial z} = Q_{net} \quad \text{on } \Gamma$$

$$\boldsymbol{\tau} = \rho_a C_\tau \llbracket \mathbf{u} \rrbracket \llbracket \mathbf{u} \rrbracket \quad Q_{net} = \mathcal{R} + \rho_a C_Q \llbracket \mathbf{u} \rrbracket \llbracket T \rrbracket$$



Velocity and temperature jump at interface

$$\llbracket \mathbf{u} \rrbracket = \mathbf{u}_a - \mathbf{u}_o \quad \text{on } \Gamma$$

$$\llbracket T \rrbracket = T_a - T_o \quad \text{on } \Gamma$$

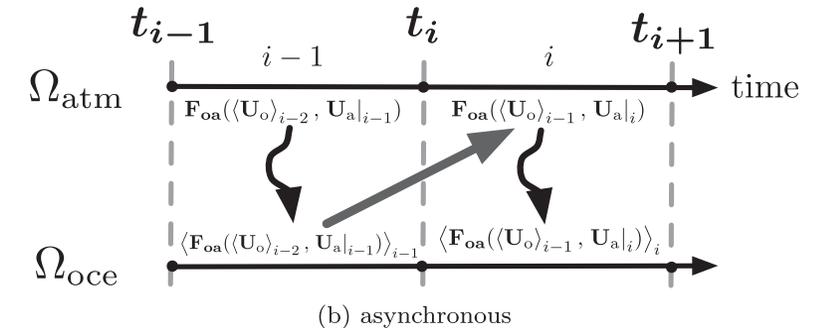
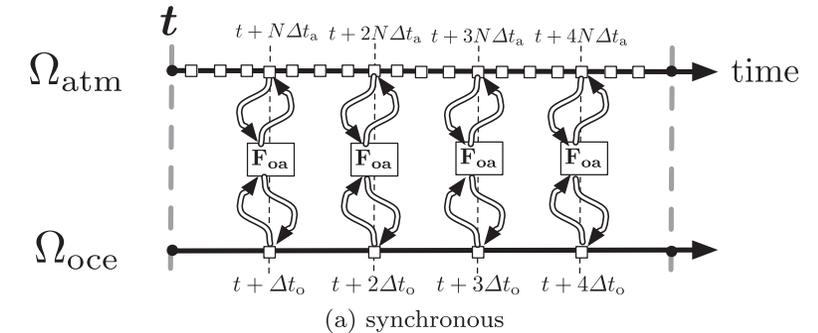
### Typical coupling methods

#### ○ Synchronous coupling

- Exchange instantaneous boundary data at largest time step
- More frequent communication
- Can be unstable

#### ○ Asynchronous coupling

- Exchange time-averaged boundary data
- Long time intervals require fewer communications between models
- Ensures flux conservation



Schematic of coupling approaches  
from Gross et al. (2018)

*Both methods can be shown to be equivalent to one step of a Schwartz algorithm*

*Recent work has investigated relationship between coupling schemes and solution methods for the monolithic ocean-atmosphere system*

- Lemarié, Blayo, Debreu (2015): Global-in-time Schwarz method
- Beljaars et al. (2017): Stable parametrized implicit flux coupling for temperature diffusion equation in the context of ice-atmosphere models
- Pelletier, Lemarié, Blayo (2017): Coupling methods for time-dependent Ekman boundary layer model
- Connors, Ganis (2011): Fluid-fluid interaction using a monolithic and a two-way partitioned method.
- Connors, Howell, Layton (2012): Partitioned methods for fluid-fluid interaction

### Our approach:

- Consider a simplified scalar equation with representative coupling conditions
- Starting from the monolithic system, develop a non-iterative approach to approximate the Neumann coupling condition
- Use a Lagrange multiplier to ensure flux continuity at the interface
- Motivated by the Implicit Value Recovery (IVR) approach applied to solid mechanics and advection-diffusion problems

Atmosphere/ocean tracer

$$\dot{T}_a + \frac{\partial}{\partial x}(u_a T_a) = \frac{\partial}{\partial z} K_a \frac{\partial T_a}{\partial z}$$

$$K_a \frac{\partial T_a}{\partial z} = K_o \frac{\partial T_o}{\partial z} = \alpha(T_a - T_o) \quad \Gamma$$

$$\dot{T}_o + \frac{\partial}{\partial x}(u_o T_o) = \frac{\partial}{\partial z} K_o \frac{\partial T_o}{\partial z}$$



## Mixed Formulation

$$\begin{aligned} \dot{\varphi}_1 - \nabla \cdot F_1(\varphi_1) &= f_1 \text{ in } \Omega_1 \\ F_1 \cdot \mathbf{n}_1 &= -\lambda \text{ on } \gamma \\ \dot{\varphi}_2 - \nabla \cdot F_2(\varphi_2) &= f_2 \text{ in } \Omega_2 \\ F_2 \cdot \mathbf{n}_2 &= \lambda \text{ on } \gamma \\ \varphi_1 &= \varphi_2 \text{ on } \gamma \end{aligned}$$

## Discretize

$$\begin{aligned} \varphi_1 &\in S_1^h \subset H_{\Gamma_1}^1(\Omega_1) \\ \varphi_2 &\in S_2^h \subset H_{\Gamma_2}^1(\Omega_2) \\ \lambda &\in G_\gamma^h \subset H^{-1/2}(\gamma) \end{aligned}$$

Semi-Discrete System  
Index 2 DAE

$$\begin{aligned} M_1 \dot{\varphi}_1 + G_1^T \lambda &= \mathbf{f}_1(\varphi_1) \\ M_2 \dot{\varphi}_2 - G_2^T \lambda &= \mathbf{f}_2(\varphi_2) \\ G_1 \varphi_1 - G_2 \varphi_2 &= 0 \end{aligned}$$

Conversion to  
Index 1 DAE

$$\begin{aligned} M_1 \dot{\varphi}_1 + G_1^T \lambda &= \mathbf{f}_1(\varphi_1) \\ M_2 \dot{\varphi}_2 - G_2^T \lambda &= \mathbf{f}_2(\varphi_2) \\ G_1 \dot{\varphi}_1 - G_2 \dot{\varphi}_2 &= 0 \end{aligned}$$

## Algebraic Form

$$\begin{bmatrix} M_1 & 0 & G_1^T \\ 0 & M_2 & -G_2^T \\ G_1 & -G_2 & 0 \end{bmatrix} \begin{bmatrix} \dot{\varphi}_1 \\ \dot{\varphi}_2 \\ \lambda \end{bmatrix} = \begin{bmatrix} \mathbf{f}_1(\varphi_1) \\ \mathbf{f}_2(\varphi_2) \\ 0 \end{bmatrix}$$

$$\begin{aligned} \text{Mass matrix } (M_i)_{kl} &= (N_{i,k}, N_{i,l})_\Omega \\ \text{Coupling matrix } (G_i)_{kl} &= (N_{i,k}, \nu_l)_\gamma \\ \text{Force vector } \mathbf{f}_{i,k} &= -(\nabla N_{i,k}, F_i)_\Omega + (N_{i,k}, f_i)_\Omega \end{aligned}$$

- Defines  $\lambda$  as an implicit function of states: can solve for  $\lambda$  and use as Neumann data
- Explicit time integration effectively decouples the subdomain equations
- Inf-sup condition verified for mortar elements
- No splitting error or stability issues



## Mixed Formulation

$$\begin{aligned} \dot{\varphi}_1 - \nabla \cdot F_1(\varphi_1) &= f_1 \text{ in } \Omega_1 \\ F_1 \cdot \mathbf{n}_1 &= -\lambda \text{ on } \gamma \\ \dot{\varphi}_2 - \nabla \cdot F_2(\varphi_2) &= f_2 \text{ in } \Omega_2 \\ F_2 \cdot \mathbf{n}_2 &= \lambda \text{ on } \gamma \\ \varphi_1 &= \varphi_2 \text{ on } \gamma \end{aligned}$$

## Discretize

$$\begin{aligned} \varphi_1 &\in S_1^h \subset H_{\Gamma_1}^1(\Omega_1) \\ \varphi_2 &\in S_2^h \subset H_{\Gamma_2}^1(\Omega_2) \\ \lambda &\in G_\gamma^h \subset H^{-1/2}(\gamma) \end{aligned}$$

Semi-Discrete System  
Index 2 DAE

$$\begin{aligned} M_1 \dot{\varphi}_1 + G_1^T \lambda &= \mathbf{f}_1(\varphi_1) \\ M_2 \dot{\varphi}_2 - G_2^T \lambda &= \mathbf{f}_2(\varphi_2) \\ G_1 \varphi_1 - G_2 \varphi_2 &= 0 \end{aligned}$$

Conversion to  
index 1 DAE

$$\begin{aligned} M_1 \dot{\varphi}_1 + G_1^T \lambda &= \mathbf{f}_1(\varphi_1) \\ M_2 \dot{\varphi}_2 - G_2^T \lambda &= \mathbf{f}_2(\varphi_2) \\ G_1 \dot{\varphi}_1 - G_2 \dot{\varphi}_2 &= 0 \end{aligned}$$

## Algebraic Form

$$\begin{bmatrix} M_1 & 0 & G_1^T \\ 0 & M_2 & -G_2^T \\ G_1 & -G_2 & 0 \end{bmatrix} \begin{bmatrix} \dot{\varphi}_1 \\ \dot{\varphi}_2 \\ \lambda \end{bmatrix} = \begin{bmatrix} \mathbf{f}_1(\varphi_1) \\ \mathbf{f}_2(\varphi_2) \\ 0 \end{bmatrix}$$

*Want to derive a similar scheme  
for bulk conditions on interface:*

$$F_1 \cdot \mathbf{n}_1 = -F_2 \cdot \mathbf{n}_2 = \alpha(\varphi_1 - \varphi_2)$$

## BULK IMPLICIT VALUE RECOVERY



Start with monolithic mixed-like formulation obtained by introducing a new flux variable  $\lambda$  and adding the bulk condition as a third equation

$$\begin{aligned} \dot{T}_a + \frac{\partial}{\partial x}(u_a T_a) &= \frac{\partial}{\partial z} K_a \frac{\partial T_a}{\partial z} & \text{in } \Omega_a & \quad \dot{T}_o + \frac{\partial}{\partial x}(u_o T_o) &= \frac{\partial}{\partial z} K_o \frac{\partial T_o}{\partial z} & \text{in } \Omega_o \\ K_a \frac{\partial T_a}{\partial z} &= \lambda & \text{on } \Gamma & \quad K_o \frac{\partial T_o}{\partial z} &= -\lambda & \text{on } \Gamma \\ \lambda &= \alpha(T_a - T_o) & \text{on } \Gamma & \end{aligned}$$

Discretize in Space: Seek  $\{T_a^h, T_o^h, \lambda^h\} \in S_{a,\Gamma}^h(\Omega_a) \times S_{o,\Gamma}^h(\Omega_o) \times G_\Gamma^h$

$$\begin{aligned} (\dot{T}_a, \psi_a)_{0,\Omega_a} + \langle \lambda, \psi_a \rangle_\Gamma &= (f_a, \psi_a)_{0,\Omega_a} + \left( K_a \frac{\partial T_a}{\partial z}, \frac{\partial \psi_a}{\partial z} \right)_{0,\Omega_a} - \left( u_a \frac{\partial T_a}{\partial x}, \psi_a \right)_{0,\Omega_a} & \forall \psi_a \in H_\Gamma^1(\Omega_a) \\ (\dot{T}_o, \psi_o)_{0,\Omega_o} - \langle \lambda, \psi_o \rangle_\Gamma &= (f_o, \psi_o)_{0,\Omega_o} + \left( K_o \frac{\partial T_o}{\partial z}, \frac{\partial \psi_o}{\partial z} \right)_{0,\Omega_o} - \left( u_o \frac{\partial T_o}{\partial x}, \psi_o \right)_{0,\Omega_o} & \forall \psi_o \in H_\Gamma^1(\Omega_o) \\ \langle \alpha(T_a - T_o) - \lambda, \mu \rangle_\Gamma dS &= 0 & \forall \mu \in H^{-1/2}(\Gamma) \end{aligned}$$

Weak form of the additional bulk condition equation

## Semi-discrete System

$$\begin{aligned}M_a \dot{\mathbf{T}}_a + G_a^T \boldsymbol{\lambda} &= \mathbf{f}_a(\mathbf{T}_a) \\M_o \dot{\mathbf{T}}_o - G_o^T \boldsymbol{\lambda} &= \mathbf{f}_o(\mathbf{T}_o) \\ \alpha G_a \mathbf{T}_a - \alpha G_o \mathbf{T}_o - \widehat{M}_\Gamma \boldsymbol{\lambda} &= 0\end{aligned}$$

Mass matrix  $(M_i)_{kl} = (N_{i,k}, N_{i,l})_\Omega$

Coupling matrix  $(G_i)_{kl} = (N_{i,k}, \nu_l)_\Gamma$

Interface mass matrix  $(\widehat{M}_\Gamma)_{kl} = (\nu_k, \nu_l)_\Gamma$

## Semi-discrete System

$$\begin{aligned}
 M_a \dot{\mathbf{T}}_a + G_a^T \boldsymbol{\lambda} &= \mathbf{f}_a(\mathbf{T}_a) \\
 M_o \dot{\mathbf{T}}_o - G_o^T \boldsymbol{\lambda} &= \mathbf{f}_o(\mathbf{T}_o) \\
 \alpha G_a \mathbf{T}_a - \alpha G_o \mathbf{T}_o - \widehat{M}_\Gamma \boldsymbol{\lambda} &= 0
 \end{aligned}$$

$$\text{Mass matrix } (M_i)_{kl} = (N_{i,k}, N_{i,l})_\Omega$$

$$\text{Coupling matrix } (G_i)_{kl} = (N_{i,k}, \nu_l)_\Gamma$$

$$\text{Interface mass matrix } (\widehat{M}_\Gamma)_{kl} = (\nu_k, \nu_l)_\Gamma$$

*Similar in form to IVR system, but cannot simplify by using time derivative of solution on interface.*

## Semi-discrete System

$$\begin{aligned}
 M_a \dot{\mathbf{T}}_a + G_a^T \boldsymbol{\lambda} &= \mathbf{f}_a(\mathbf{T}_a) \\
 M_o \dot{\mathbf{T}}_o - G_o^T \boldsymbol{\lambda} &= \mathbf{f}_o(\mathbf{T}_o) \\
 \alpha G_a \mathbf{T}_a - \alpha G_o \mathbf{T}_o - \widehat{M}_\Gamma \boldsymbol{\lambda} &= 0
 \end{aligned}$$

Mass matrix  $(M_i)_{kl} = (N_{i,k}, N_{i,l})_\Omega$

Coupling matrix  $(G_i)_{kl} = (N_{i,k}, \nu_l)_\Gamma$

Interface mass matrix  $(\widehat{M}_\Gamma)_{kl} = (\nu_k, \nu_l)_\Gamma$

*Similar in form to IVR system, but cannot simplify by using time derivative of solution on interface.*

**Solution:** Discretize in time, then solve the fully discrete problem for flux  $\boldsymbol{\lambda}$

$$\begin{aligned}
 M_a \left( \frac{\mathbf{T}_a^{n+1} - \mathbf{T}_a^n}{\Delta t} \right) + G_a^T \boldsymbol{\lambda} &= \mathbf{f}_a(\mathbf{T}_a^n) \\
 M_o \left( \frac{\mathbf{T}_o^{n+1} - \mathbf{T}_o^n}{\Delta t} \right) - G_o^T \boldsymbol{\lambda} &= \mathbf{f}_o(\mathbf{T}_o^n) \\
 \alpha G_a \mathbf{T}_a^{n+1} - \alpha G_o \mathbf{T}_o^{n+1} - \widehat{M}_\Gamma \boldsymbol{\lambda} &= 0
 \end{aligned}$$

Separate system into internal ( $I$ ) and interface ( $\Gamma$ ) degrees of freedom

$$\mathbf{g}_i(\mathbf{T}_i^n) = \Delta t \mathbf{f}_i(\mathbf{T}_i^n) - M_i \mathbf{T}_i^n$$

$$\begin{bmatrix} M_{a,\Gamma\Gamma} & 0 & \Delta t G_a^T & M_{a,\Gamma I} & 0 \\ 0 & M_{o,\Gamma\Gamma} & -\Delta t G_o^T & 0 & M_{o,\Gamma I} \\ \alpha G_a & -\alpha G_o & -\widehat{M}_\Gamma & 0 & 0 \\ \hline M_{a,I\Gamma} & 0 & 0 & M_{a,II} & 0 \\ 0 & M_{o,I\Gamma} & 0 & 0 & M_{o,II} \end{bmatrix} \begin{bmatrix} \mathbf{T}_{a,\Gamma}^{n+1} \\ \mathbf{T}_{o,\Gamma}^{n+1} \\ \boldsymbol{\lambda} \\ \mathbf{T}_{a,I}^{n+1} \\ \mathbf{T}_{o,I}^{n+1} \end{bmatrix} = \begin{bmatrix} \mathbf{g}_{a,\Gamma}(\mathbf{T}_a^n) \\ \mathbf{g}_{o,\Gamma}(\mathbf{T}_o^n) \\ 0 \\ \mathbf{g}_{a,I}(\mathbf{T}_a^n) \\ \mathbf{g}_{o,I}(\mathbf{T}_o^n) \end{bmatrix}$$

Solve for flux: **with explicit time stepping only involves information from old time step!**

$$\boldsymbol{\lambda} = \left( \Delta t G_a^T A_a^{-1} G_a + \Delta t G_o^T A_o^{-1} G_o - \frac{\widehat{M}_\Gamma}{\alpha} \right)^{-1} \left( G_a^T A_a^{-1} \hat{\mathbf{g}}_a(\mathbf{T}_a^n) - G_o^T A_o^{-1} \hat{\mathbf{g}}_o(\mathbf{T}_o^n) \right)$$

where  $\hat{\mathbf{g}}_i(\mathbf{T}_i^n) = \mathbf{g}_{i,\Gamma}(\mathbf{T}_i^n) - M_{i,\Gamma I} M_{i,II}^{-1} \mathbf{g}_{i,I}(\mathbf{T}_i^n)$

$$A_i = M_{i,\Gamma} - M_{i,\Gamma I} M_{i,II}^{-1} M_{i,I\Gamma}$$

1. Compute right-hand side terms

$$\mathbf{g}_i(\mathbf{T}_i^n) = \Delta t \mathbf{f}_i(\mathbf{T}_i^n) - M_i \mathbf{T}_i^n$$

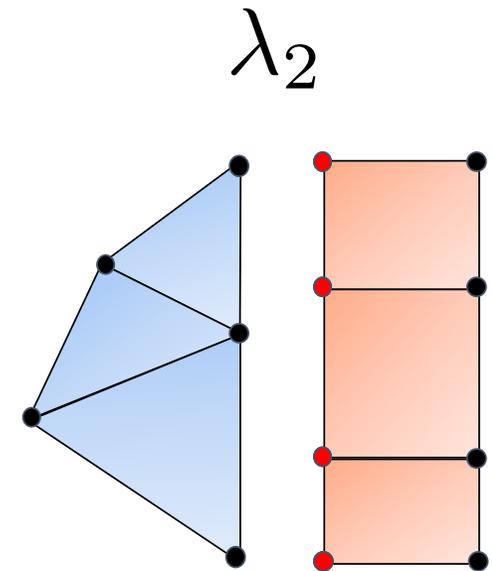
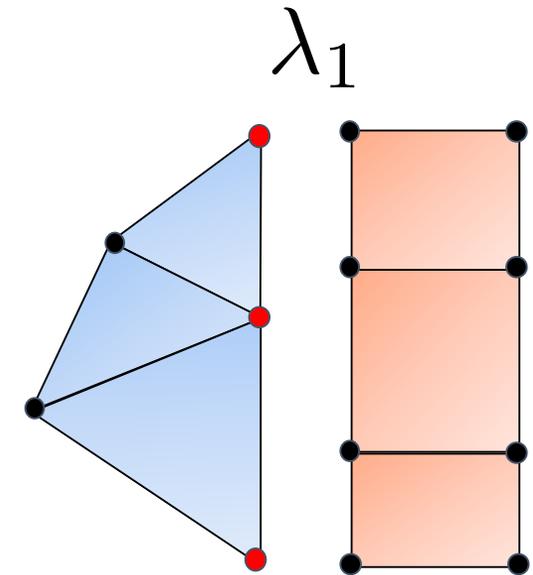
2. Estimate interface boundary condition

$$\boldsymbol{\lambda} = \left( \Delta t G_a^T A_a^{-1} G_a + \Delta t G_o^T A_o^{-1} G_o - \frac{\widehat{M}_\Gamma}{\alpha} \right)^{-1} (G_a^T A_a^{-1} \hat{\mathbf{g}}_a(\mathbf{T}_a^n) - G_o^T A_o^{-1} \hat{\mathbf{g}}_o(\mathbf{T}_o^n))$$

3. Solve independently in each subdomain

$$\begin{bmatrix} M_{i,\Gamma} & M_{i,\Gamma I} \\ M_{i,I\Gamma} & M_{i,II} \end{bmatrix} \begin{bmatrix} \mathbf{T}_{i,\Gamma} \\ \mathbf{T}_{i,I} \end{bmatrix} = \begin{bmatrix} \mathbf{g}_{i,\Gamma}^n \pm G_i^T \boldsymbol{\lambda} \\ \mathbf{g}_{i,I}^n \end{bmatrix}$$

- There is some flexibility in choosing the Lagrange multiplier space
- For the original IVR formulation, we followed the mortar method approach and chose either one of the interface partitions
- Results in a formulation that satisfies the inf-sup condition
- We follow this approach in the bulk IVR method
- Expect to converge optimally, but not pass a patch test



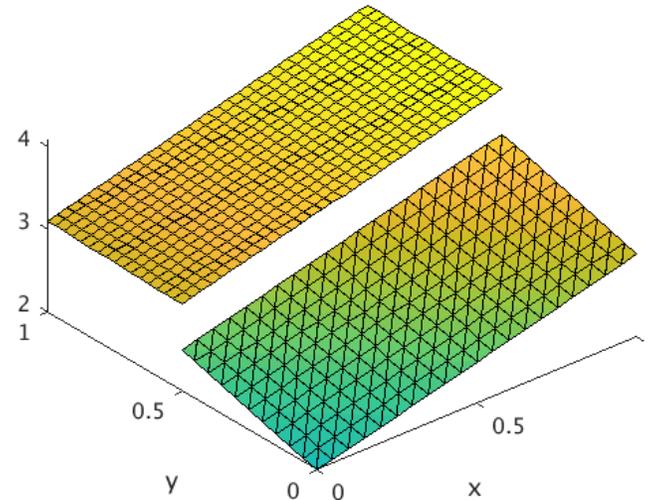
$$T_o = x + y + 2$$

$$T_a = x + \frac{y}{10} + 3$$

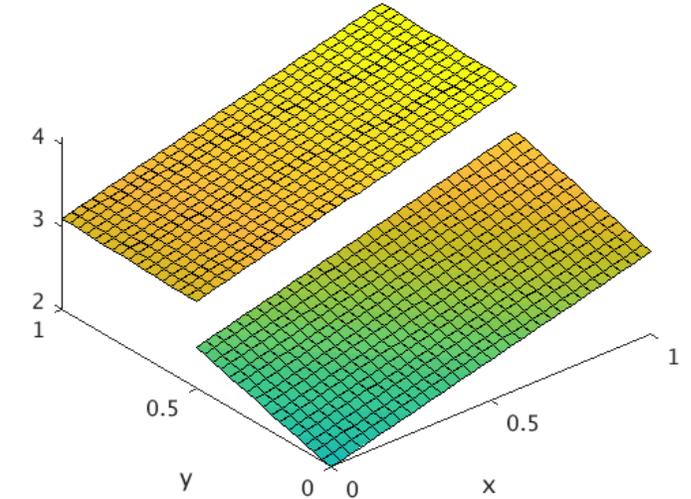
$$K_o = 0.001$$

$$K_a = 0.01$$

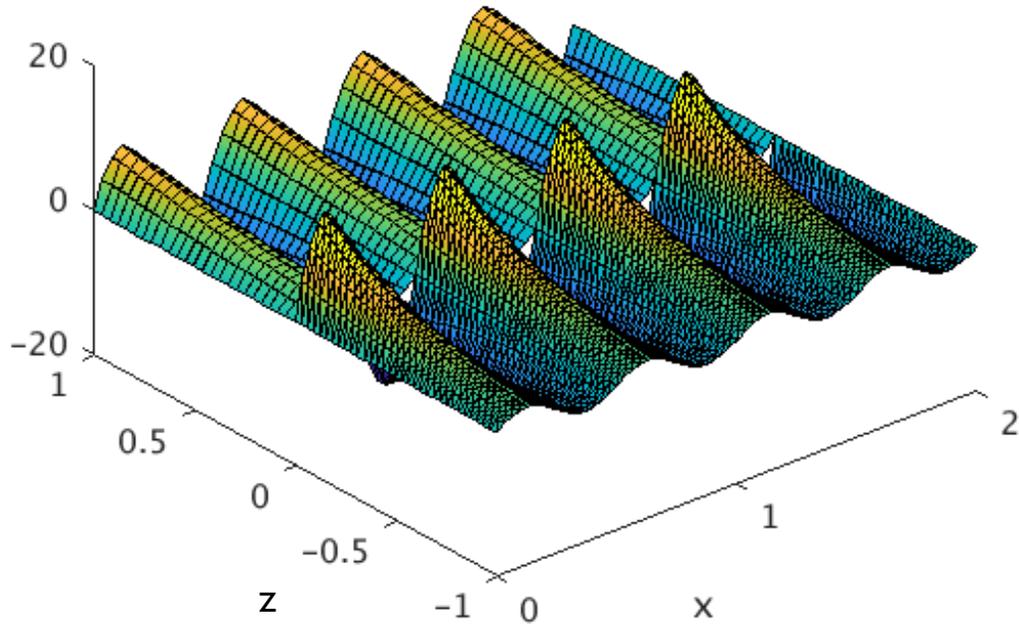
Non-matching grids



Matching grids



Error Norm	$h(\Omega_o)$	$h(\Omega_a)$	BIVR ( $\lambda_o$ )	BIVR ( $\lambda_a$ )
$L^2(\Omega)$	0.03125	0.03125	1.23e-15	1.23e-15
$H^1(\Omega)$	0.03125	0.03125	1.26e-13	1.26e-13
$L^2(\Omega)$	0.05000	0.03125	7.18e-06	2.14e-07
$H^1(\Omega)$	0.05000	0.03125	7.77e-04	1.64e-05



$$T_o = d_o \exp(\beta_o z) \sin(n\pi x) \exp(\omega t)$$

$$T_a = d_a \exp(\beta_a z) \sin(n\pi x) \exp(\omega t)$$

$$\beta_o = 2$$

$$\beta_a = 1$$

$$d_o = 5$$

$$d_a = 1$$

$$K_o = 0.001$$

$$K_a = 0.01$$

$$\alpha = \frac{K_o \beta_o d_o}{d_a - d_o}$$

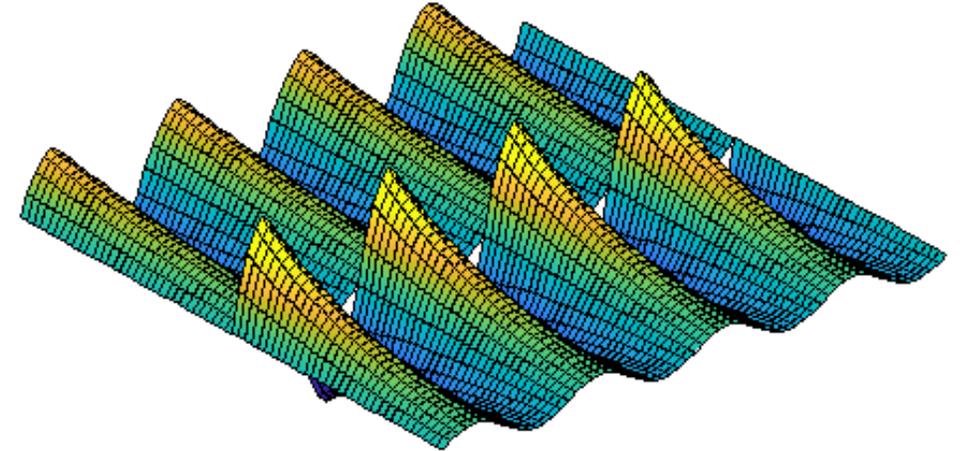
$$u_o = u_a = 1$$

Note: developed for testing method convergence and not intended to be physically realistic example.

$$T_o = d_o \exp(\beta_o z) \sin(n\pi x) \exp(\omega t)$$

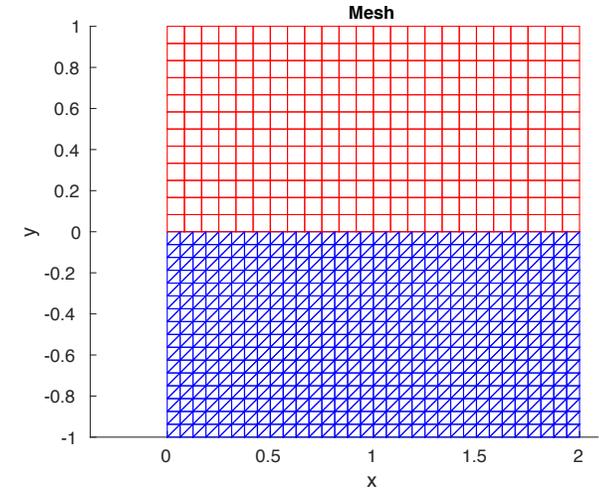
$$T_a = d_a \exp(\beta_a z) \sin(n\pi x) \exp(\omega t)$$

Matching grid solution

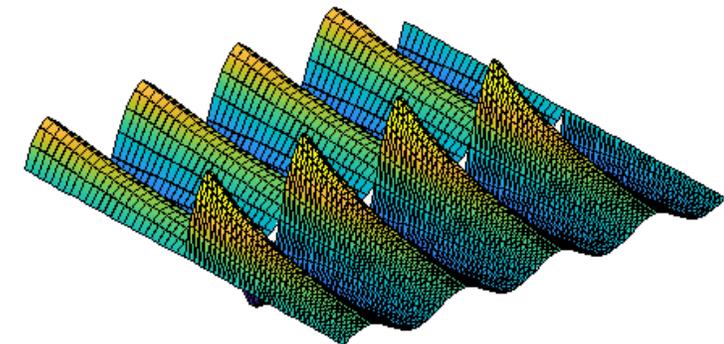


Mesh ( $\Omega_o$ )	Mesh ( $\Omega_a$ )	$\Delta t$	$L^2(\Omega)$	$H^1(\Omega)$
$16 \times 8$	$16 \times 8$	1.89e-02	1.44e-00	4.86e01
$32 \times 16$	$32 \times 16$	9.43e-03	2.50e-01	2.38e01
$64 \times 32$	$64 \times 32$	4.69e-03	4.55e-02	1.19e01
$128 \times 64$	$128 \times 64$	1.83e-03	8.76e-03	5.92e00
Rate	-	-	<b>2.38</b>	<b>1.01</b>

Non-matching grids



Solution

 $L^2(\Omega)$  Error norm

Mesh ( $\Omega_o$ )	Mesh ( $\Omega_a$ )	$\Delta t$	BIVR( $\lambda_o$ )	BIVR( $\lambda_a$ )
$16 \times 8$	$12 \times 6$	1.33e-02	2.09e-00	2.09e-00
$32 \times 16$	$24 \times 12$	6.67e-03	3.40e-01	3.40e-01
$64 \times 32$	$48 \times 24$	3.32e-03	6.18e-02	6.18e-02
$128 \times 64$	$96 \times 48$	1.66e-03	1.30e-02	1.30e-02
Rate	-	-	2.25	2.25

 $H^1(\Omega)$  Error norm

Mesh ( $\Omega_o$ )	Mesh ( $\Omega_a$ )	$\Delta t$	BIVR( $\lambda_o$ )	BIVR( $\lambda_a$ )
$16 \times 8$	$12 \times 6$	1.33e-02	5.66e01	5.66e01
$32 \times 16$	$24 \times 12$	6.67e-03	2.78e01	2.78e01
$64 \times 32$	$48 \times 24$	3.32e-03	1.37e01	1.37e01
$128 \times 64$	$96 \times 48$	1.66e-03	6.84e00	6.84e00
Rate	-	-	1.01	1.01

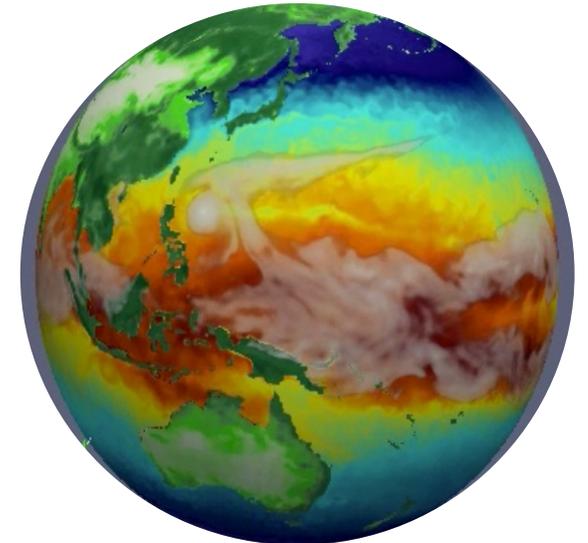
*Extended IVR to a Bulk-IVR partitioned scheme for a scalar equation with bulk coupling conditions*

- Starts with a well-posed monolithic mixed-like formulation
- Explicit time integration results in an IVR-like structure
- This structure enables solving for the flux on the interface
- Results in a non-iterative partitioned scheme
- Proof-of-concept tested on simple manufactured solutions

*Next steps*

- Extend Bulk-IVR to simplified coupled fluid equations
- Extend to conjugate heat transfer with imperfect transmission conditions
- Investigate extensions to non-linear coupling conditions
- Evaluate accuracy and stability of method for different spatial and time discretizations

$$\begin{aligned} \dot{T}_a + \frac{\partial}{\partial x}(u_a T_a) &= \frac{\partial}{\partial z} K_a \frac{\partial T_a}{\partial z} \\ \updownarrow & \boxed{K_a \frac{\partial T_a}{\partial z} = K_o \frac{\partial T_o}{\partial z} = \alpha(T_a - T_o)} \Gamma \\ \dot{T}_o + \frac{\partial}{\partial x}(u_o T_o) &= \frac{\partial}{\partial z} K_o \frac{\partial T_o}{\partial z} \end{aligned}$$





A. Beljaars, E. Dutra, G. Balsama, F. Lemarie (2017), On the numerical stability of surface-atmosphere coupling in weather and climate models, *Geosci. Model Dev.*, 10:977-989.

J. Connors, B. Ganis (2011), Stability of algorithms for a two domain natural convection problem and observed model uncertainty, *Comput. Geosci.*, 15:509-527.

J. Connors, J. Howell, and W. Layton (2012), Decoupled time stepping methods for fluid-fluid interaction, *SINUM*, 50, No. 3, 1297-1319.

M. Gross, *et al.* (2018) Physics-Dynamics Coupling in Weather, Climate, and Earth System Models: Challenges and Recent Progress, *MWR*, 146:3505-3544.

F. Lemarié, E. Blayo, and L. Debreu, (2015) Analysis of ocean-atmosphere coupling algorithms: Consistency and stability, *Proc. Comput. Sci.*, 51, 2066-2075.

C. Pelletier, F. Lemarié, and E. Blayo, (2017) A Theoretical study of a simplified air-sea coupling problem including turbulent parameterizations, *Coupled Problems 2017*, 38-48.

K. Peterson, P. Bochev, P. Kubbery, (2018) Explicit non-iterative partitioned algorithms for interface problems based on Lagrange multipliers, *CAMWA*, doi:10.1016/j.camwa.2018.09.045.